# Theory and Practice For Calibration of Moving Coil Type Seismometer. 

2nd Edition

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## 1. Introduction

The equation of motion of the pendulum mass for moving coil type seismometer (Fig. 1) is
$\ddot{x}(t)+2 \omega_{0} h \dot{x}(t)+\omega_{0}{ }^{2} x(t)=-\ddot{y}(t)$,
where $x(t)$ denotes the displacement of the pendulum mass relative to the flame of seismometer, $\ddot{y}(t)$ the input ground acceleration, $\omega_{0}$ the natural angular frequency ( $\omega_{0}=2 \pi f_{0}=2 \pi / T_{0}$ ), $f_{0}$ the natural frequency, $T_{0}$ the natural period, $h$ the damping constant which is the sum of the mechanical one $h_{m}$ and electro-magnetic one $h_{e}$ which is controlled by the shunt resistance $R_{s}$ (e.g., Yokoi (1996)). The discussion here is always based on this equation of motion.

If we can apply a controlled input ground displacement $y(t)$ or acceleration $\ddot{y}(t)$, direct calibration can be performed. Controlled $y(t)$, e.g., can be provided by a shaking table.

Calibration can be applied also by making use of the electro-magnetic force induced by the input current that is given to the transducer coil itself or to the calibration coil if exists. Willmore impedancebridge is an effective tool to separate the output current from input one in case of single coil seismometer (e.g., Willmore, 1959).


Fig. 1 Simplified sketch of electro-magnetic moving coil seismograph. When a coil attached to the mass moves through the magnetic field, the voltage across the coil terminal is proportional to the relative velocity between mass and magnet. Redrawn based on Aki and Richards (1980).

## 2. Theory and Method

### 2.1. Calibration on Shaking Table

The output voltage of the seismometer directly connected to the parallel shunt resistance $R_{s}$ as shown in Fig. 2 is given as follows.

$$
\begin{equation*}
e(t)=\frac{G R_{s}}{R_{0}+R_{s}} \cdot \dot{x}(t) \tag{2}
\end{equation*}
$$

where $G$ is the electro-dynamical constant so called the sensitivity, $R_{0}$ the coil resistance of the seismometer. A shaking table gives a sinusoidal input displacement

$$
y(t)=y_{m} \exp (i \omega t)
$$

with arbitrary angular frequency $\omega=2 \pi f$, where $f$ is the frequency. The mass displacement and the output voltage may also have sinusoidal form.

$$
\begin{aligned}
& x(t)=x_{m} \exp (i \omega t) \\
& e(t)=e_{m} \exp (i \omega t)
\end{aligned}
$$

Then, Eq.(1) can be solved in the form

$$
\begin{equation*}
\frac{-x_{m}}{y_{m}}=\frac{1}{1-2 i h\left(\omega_{0} / \omega\right)-\left(\omega_{0} / \omega\right)^{2}} . \tag{3}
\end{equation*}
$$

The ratio of the output voltage $e_{m}$ to the input ground velocity $i \omega y_{m}$ at $\omega$ is derived from Eq. (2) and (3).

$$
\begin{align*}
\frac{e_{m}}{i \omega y_{m}} & =\frac{-e_{m}}{i \omega x_{m}} \cdot \frac{-i \omega x_{m}}{i \omega y_{m}} \\
& =\frac{G R_{s}}{R_{0}+R_{s}} \cdot \frac{1}{1-2 i h\left(\omega_{0} / \omega\right)-\left(\omega_{0} / \omega\right)^{2}} \tag{4}
\end{align*}
$$

Consider the relation $\omega_{0}=2 \pi f_{0}=2 \pi / T_{0}, \omega=2 \pi f=2 \pi / T$, then,

$$
\begin{equation*}
-i\left\{1+\left(R_{0}+R_{s}\right)\right\} \cdot \frac{e_{m} T}{2 \pi y_{m}}=\frac{G}{1-2 i h\left(T / T_{0}\right)-\left(T / T_{0}\right)^{2}} \tag{5}
\end{equation*}
$$

Define the function $F_{l}(T)$ by the absolute value of this function.

$$
\begin{equation*}
F_{1}(T)=\left\{1+\left(R_{0} / R_{s}\right)\right\} \cdot \frac{e_{m} T}{2 \pi y_{m}}=\frac{G}{\sqrt{\left\{1-\left(T / T_{0}\right)^{2}\right\}^{2}+4 h^{2}\left(T / T_{0}\right)^{2}}} \tag{6.1}
\end{equation*}
$$

The high frequency asymptote of $F_{l}(T)$ gives $G$ itself (Fig.3).
The argument of the both members of Eq.(5) gives the phase lag of the output voltage from the input ground displacement $\operatorname{Arg}\left(e_{m} / y_{m}\right)$. The function $P_{l}(T)$ is defined as follows.

$$
\begin{equation*}
P_{1}(T)=\operatorname{Arg}\left(e_{m} / y_{m}\right)-\frac{\pi}{2}=-\tan ^{-1}\left[-2 h\left(T / T_{0}\right) /\left[1-\left(T / T_{0}\right)^{2}\right\}\right]+2 N \pi \tag{6.2}
\end{equation*}
$$

At the natural period $T=T_{0}, P_{l}(T)$ takes the value $\pi / 2$.


Fig. 2 Circuit for the shaking table test.


Fig. 3 Theoretical Curve for $F_{l}(T)(t o p)$ and $P_{l}(T)$ (bottom).

### 2.2. Calibration using Willmore Impedance Bridge

The circuit shown in Fig. 4 is selected for convenience in the experiment. Any voltage imbalance between $A$ and $B: V_{A B}$, can not give any influence on the potential difference between $a$ and $b: V_{a b}$, because the impedance bridge is balanced.

Suppose that the pendulum of the seismometer is locked, namely, there is not any secondary induced current in the circuit. The potential difference $V_{a b}$ due to the electromotive force $E(t)$ produces the current $i(t)$ in the resistance $R$.

$$
\begin{equation*}
i=\frac{E(t)\left(1+R_{0} / R_{s}\right)}{R\left(1+R_{0} / R_{s}\right)+R_{0}} . \tag{7}
\end{equation*}
$$

If the current in the coil resistance $R_{0}$ is $I$, that in the shunt resistance $R_{s}$ must be $i-I$ and these two resistors make the same potential difference $R_{0} I=R_{s}(i-I)$. Then,

$$
\begin{equation*}
I=\frac{i(t)}{1+R_{0} / R_{s}}=\frac{E(t)}{R\left(1+R_{0} / R_{s}\right)+R_{0}} . \tag{8}
\end{equation*}
$$

This is the primary current input given to the seismometer by the function generator.


Fig. 4 Circuit for Willmore impedance bridge (Redrawn after Aoki (1994)).
Note: (-) and (G) of the seismometer should be connected together with (-) of the recorder. (G) and (-) of the recorder should not be connected.

Release the pendulum. The current introduced to the seismometer's coil induces an electro-magnetic force that applies to the pendulum.

$$
\begin{equation*}
F_{e m}=G \cdot I=\frac{G E(t)}{R\left(1+R_{0} / R_{s}\right)+R_{0}} \tag{9}
\end{equation*}
$$

where $m$ is the weight of the pendulum mass for the linear motion type and the ratio of the inertial moment of the pendulum to the arm length of electromagnetic force for the rotation type. In the absence of input ground motion: $\ddot{y}(t) \equiv 0$.

This force must appear in the equation of motion as follows.

$$
\begin{equation*}
\ddot{x}(t)+2 \omega_{0} h \dot{x}(t)+\omega_{0}^{2} x(t)=\frac{-G E(t)}{m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}}-\ddot{y}(t) . \tag{10}
\end{equation*}
$$

The velocity of mass motion $\dot{x}(t)$ produces the secondary electromotive force

$$
E^{s}(t)=G \cdot \dot{x}(t),
$$

and the secondary induced current $I^{s}(t)$ i.e., the output of the seismometer.

$$
I^{s}(t)=\frac{E^{s}(t)}{R_{0}+R_{s}^{\prime}}=\frac{G}{R_{0}+R_{s}^{\prime}} \cdot \dot{x}(t)
$$

where $R_{s}^{\prime}$ is total shunt resistance which consists of all resistors in the impedance bridge.

$$
R_{s}^{\prime}=\frac{1}{\left(1 / R_{s}\right)+\left\{2 R+\left(1 / R_{0}+1 / R_{s}\right)^{-1}\right\}^{-1}}=R_{s} \cdot \frac{\left(1+R_{0} / R_{s}+R_{0} / 2 R\right)}{\left(1+R_{0} / R_{s}\right) \cdot\left(1+R_{s} / 2 R\right)+R_{0} / 2 R}
$$

$R_{s}^{\prime}$ tends to $R_{s}$ at $R \gg R_{s}>R_{0}$, and to $2 R+R_{0}$ at infinite $R_{s}$.
The potential difference $V_{a A}$, produced by this current can be calculated as follows

$$
\left\{2 R+\left(R_{0}^{-1}+R_{s}^{-1}\right)^{-1}\right\} \cdot I_{a b}=R_{s}\left(I^{s}-I_{a b}\right)
$$

Then, $\quad I_{a b}=I^{s} \cdot \frac{R_{s}}{R_{s}+2 R+\left(R_{0}^{-1}+R_{s}^{-1}\right)^{-1}}$
The potential difference $V_{a b}$ is

$$
V_{a b}=-2 R I_{a b}=-2 R I^{s} \cdot \frac{R_{s}}{R_{s}+2 R+\left(R_{0}^{-1}+R_{s}^{-1}\right)^{-1}}=-R_{s}^{\prime \prime} I^{s},
$$

where

$$
\begin{aligned}
\mathrm{R}_{\mathrm{s}}^{\prime \prime} & =2 \mathrm{R} \cdot \frac{\mathrm{R}_{\mathrm{s}}}{\mathrm{R}_{\mathrm{s}}+2 \mathrm{R}+\left(\mathrm{R}_{0}^{-1}+\mathrm{R}_{\mathrm{s}}^{-1}\right)^{-1}}=\mathrm{R}_{\mathrm{s}} \cdot \frac{1+\mathrm{R}_{0} / \mathrm{R}_{\mathrm{s}}}{\left(1+\mathrm{R}_{0} / \mathrm{R}_{\mathrm{s}}\right)+\left(\mathrm{R}_{\mathrm{s}}^{-1}+2 \mathrm{R}_{0}\right) / 2 \mathrm{R}} \\
& =2 \mathrm{R} \cdot \frac{1}{1+2 \mathrm{R} / \mathrm{R}_{\mathrm{s}}+\mathrm{R}_{0} /\left(\mathrm{R}_{0}+\mathrm{R}_{\mathrm{s}}\right)} .
\end{aligned}
$$

This also tends to $R_{s}$ at $R \gg R_{s}>R_{0}$, and to $2 R$ at infinite $R_{s}$
Then, $\quad V_{a b}=-\frac{G R_{s}^{\prime}}{R_{0}+R_{s}^{\prime}} \cdot \dot{x}(t)$.

### 2.2.1. Sinusoidal Current Input.

Assume the sinusoidal feature with angular frequency $\omega$ for all signals, i.e.

$$
\begin{aligned}
& \left.x(t)=x_{m} \exp -i \omega t\right), \\
& y(t)=y_{m} \exp -i \omega t, \\
& E(t)=E_{m} \exp (i \omega t), \\
& \left.V_{a b}(t)=\left(V_{a b}\right)_{m} \exp -i \omega t\right),
\end{aligned}
$$

The ratio of relative mass motion to the input voltage is derived in the following formula substituting them in Eq. (10).

$$
\begin{equation*}
\frac{x_{m}}{E_{m}}=\frac{-G}{m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}} \cdot \frac{1}{\left\{1-2 i h\left(\omega_{0} / \omega\right)-\left(\omega_{0} / \omega\right)^{2}\right\}\left(-\omega^{2}\right)} . \tag{13}
\end{equation*}
$$

Eq.(12) can be changed in the following form, because the differentiation on time corresponds to the product with $-i \omega$ in this case.

$$
\frac{\left(V_{a b}\right)_{m}}{x_{m}}=\frac{-i \omega G R^{\prime \prime}}{R_{0}+R^{\prime}} .
$$

Then, the ratio of the output voltage $\left(V_{a b}\right)_{m}$ to the input voltage imbalance $E_{m}$ is given as

$$
\begin{align*}
\frac{\left(V_{a b}\right)_{m}}{E_{m}} & =\frac{\left(V_{a b}\right)_{m}}{x_{m}} \cdot \frac{x_{m}}{E_{m}} \\
& =\frac{-i R^{\prime \prime}{ }_{s}}{R_{0}+R_{s}^{\prime}} \cdot \frac{G^{2}}{\omega m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}} \cdot \frac{1}{\left\{1-2 i h\left(\omega_{0} / \omega\right)\left(\omega_{0} / \omega\right)^{2}\right\}} . \tag{14}
\end{align*}
$$

The absolute value is

$$
\begin{equation*}
\left|\frac{\left(V_{a b}\right)_{m}}{E_{m}}\right|=\frac{R_{s}^{\prime \prime}}{R_{0}+R_{s}^{\prime}} \cdot \frac{G^{2}}{\omega m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}} \cdot \frac{1}{\sqrt{\left\{1-\left(T / T_{0}\right)^{2}\right\}^{2}+4 h^{2}\left(T / T_{0}\right)^{2}}} . \tag{15.1}
\end{equation*}
$$

The phase lag is

$$
\begin{equation*}
\operatorname{Arg}\left(\frac{\left(V_{a b}\right)_{m}}{E_{m}}\right)=-\frac{\pi}{2}-\tan ^{-1}\left[-2 h\left(T / T_{0}\right) / 1-\left(T / T_{0}\right)^{2}\right]+2 N \pi . \tag{15.2}
\end{equation*}
$$

Define two functions $F_{2}(T)$ and $P_{2}=(T)$ as follows.

$$
\begin{equation*}
F_{2}(T)=\frac{\omega\left(R_{0}+R_{s}^{\prime}\right)\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}}{R_{s}^{\prime}} \cdot\left|\frac{\left(V_{a b}\right)_{m}}{E_{m}}\right|=\frac{G^{2}}{m} \cdot \frac{1}{\sqrt{\left\{1-\left(T / T_{0}\right)^{2}\right\}^{2}+4 h^{2}\left(T / T_{0}\right)^{2}}} . \tag{16.1}
\end{equation*}
$$

where $\omega=2 \pi / T$. The high frequency asymptote of $F_{2}(T)$ gives the value of $G^{2} / m$.

$$
\begin{align*}
P_{2}(T) & =\operatorname{Arg}\left(\frac{\left(V_{a b}\right)_{m}}{E_{m}}\right)+\frac{\pi}{2}  \tag{16.2}\\
& =-\tan ^{-1}\left[-2 h\left(T / T_{0}\right) /\left[1-\left(T / T_{0}\right)^{2}\right\}\right]+2 N \pi
\end{align*}
$$

Note that $F_{2}(T)$ is almost same as $F_{1}(T)$ except the scaling factor, and $P_{2}(T)$ is the same function as $P_{1}(T)$. Then, the shape of $F_{2}(T)$ and $P_{2}(T)$ also can be seen in Fig.3.

### 2.2.2. Step Current Input.

Assume that the input voltage imbalance is a step function $E(t)=E_{0} \cdot H(t)$, where $H(t)$ denotes the Heviside function. The time domain solution of the equation of motion

$$
\begin{equation*}
\ddot{x}(t)+2 \omega_{0} h \dot{x}(t)+\omega_{0}{ }^{2} x(t)=\frac{-G E_{0} H(t)}{m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}}, \tag{17}
\end{equation*}
$$

is given by Kitsunezaki and Goto (1964) as

$$
\dot{x}(t)=\left\{\begin{array}{lr}
-G E_{0} / m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\} \cdot e^{-\omega_{0} h t} / \omega_{0} \sqrt{1-h^{2}} \cdot \sin \left(\omega_{0} \sqrt{1-h^{2}} \cdot t\right), h<1.0,  \tag{18}\\
-G E_{0} / m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\} \cdot t \cdot e^{-\omega_{0} t}, \\
-G E_{0} / m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\} \cdot \sinh \left(\omega_{0} \sqrt{h^{2}-1} \cdot t\right), & h=1.0, \\
& h>1.0 .
\end{array}\right.
$$

Substitute it in Eq.(12) and the output voltage change is given as follows.

$$
V_{a b}(t)= \begin{cases}\left(G^{2} E_{0} R_{s}^{\prime \prime}\right) /\left[m\left(R_{0}+R_{s}^{\prime}\right)\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}\right] \cdot e^{-\omega_{0} h t} / \omega_{0} \sqrt{1-h^{2}} \cdot \sin \left(\omega_{0} \sqrt{1-h^{2}} \cdot t\right), h<1.0,  \tag{19}\\ \left(G^{2} E_{0} R_{s}^{\prime \prime}\right) /\left[m\left(R_{0}+R_{s}^{\prime}\right)\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}\right] \cdot t \cdot e^{-\omega_{0} t}, & h=1.0, \\ \left(G^{2} E_{0} R_{s}^{\prime \prime}\right) /\left[m\left(R_{0}+R_{s}^{\prime}\right)\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}\right] \cdot \sinh \left(\omega_{0} \sqrt{1-h^{2}} \cdot t\right), & h>1.0 .\end{cases}
$$

This shows the apparent period $T_{d}$, related to the natural one.

$$
T_{d}= \begin{cases}T_{0} / \sqrt{1-h^{2}}, & h<1.0, \\ T_{0}, & h=1.0, \\ T_{0} / \sqrt{h^{2}-1} & h>1.0\end{cases}
$$

The damping constant can be estimated by measuring the damping ratio, i.e., the amplitude ratio of peaks and following troughs for $h<1$. When the damping constant $h$ is close to 1.0 , it is difficult to determine its value precisely by means of damping ratio. Kitsunezaki and Goto(1964) has given the relation of $h$ with $T_{w}$, which is the time when record reduces its value to 0.2 times the maximum value.

Eq.(19) gives the amplitude of the first peak as follows,

$$
\begin{equation*}
\left(V_{a b}\right)_{f . p .}=\frac{\left(G^{2} E_{0} R_{s}^{\prime \prime}\right)}{\left[m\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}\left(R_{0}+R_{s}^{\prime}\right)\right] \omega_{0} P(h)}, \tag{20}
\end{equation*}
$$

where

$$
P(h)=\left\{\begin{array}{l}
\exp \left\{h / \sqrt{1-h^{2}} \tan ^{-1} \sqrt{1-h^{2}} / h\right\} h<1.0, \\
2.718, \\
\exp \left\{h / \sqrt{h^{2}-1} \tanh ^{-1} \sqrt{h^{2}-1} / h\right\} h>1.0 .
\end{array}\right.
$$

and the suffix ${ }_{f . p}$. denotes the amplitude of the first peak. The value of $G^{2} / m$ is given by the absolute value of Eq.(20) in the following formula.

$$
\begin{equation*}
\frac{G^{2}}{m}=\frac{\left[\left\{R\left(1+R_{0} / R_{s}\right)+R_{0}\right\}\left(R_{0}+R_{s}^{\prime}\right)\right] \omega_{0} P(h)}{E_{0} R_{s} "} \cdot\left|\left(V_{a b}\right)_{f . p .}\right| . \tag{21}
\end{equation*}
$$

### 2.3. Calibration using Calibration Coil

If the seismometer has the second coil for calibration as show in Fig.5, the output current from the transducer coil can't be mixed with the input current given to the calibration coil. Then, any impedance bridge is not required. Electro motive force given to the second coil $E_{2}$ produces the input current $I_{2}$ in the second coil.

$$
I_{2}=E_{2} / R_{2}
$$

where $R_{2}$ is the resistance of the calibration coil.
This current gives the electro-magnetic force

$$
\begin{equation*}
F_{e m}=G_{2} \cdot I_{2}, \tag{23}
\end{equation*}
$$

where $G_{2}$ is the electro-dynamical constant of the calibration coil.
The equation of motion for the pendulum mass is

$$
\begin{equation*}
\ddot{x}(t)+2 \omega_{0} h \dot{x}(t)+\omega_{0}^{2} x(t)=\frac{-G_{2} E_{2}(t)}{m R_{2}} . \tag{24}
\end{equation*}
$$



Fig. 5 Circuit for the calibration coil.

### 2.3.1. Sinusoidal Current Input

Assume the sinusoidal feature with angular frequency $\omega$ for all signals. The equation of motion is solved in the form

$$
\begin{equation*}
\frac{x_{m}}{\left(E_{2}\right)_{m}}=\frac{G_{2}}{m R_{2}} \cdot \frac{1}{\left\{1-2 i h\left(\omega_{0} / \omega\right)-\left(\omega_{0} / \omega\right)^{2}\right\} \omega^{2}} . \tag{25}
\end{equation*}
$$

The ratio of output voltage to the mass displacement is given as

$$
\frac{\left(V_{c d}\right)_{m}}{x_{m}}=\frac{-i \omega G_{1} R_{s}}{R_{0}+R_{s}}
$$

Then,

$$
\begin{equation*}
\frac{\left(V_{c d}\right)_{m}}{\left(E_{2}\right)_{m}}=\frac{\left(V_{c d}\right)_{m}}{x_{m}} \cdot \frac{x_{m}}{\left(E_{2}\right)_{m}}=\frac{-i \omega G_{1} G_{2} R_{s}}{\left(R_{0}+R_{2}\right) m R_{2}} \cdot \frac{1}{\left\{1-2 i h\left(\omega_{0} / \omega\right)\left(\omega_{0} / \omega\right)^{2}\right\} \omega^{2}} \tag{26}
\end{equation*}
$$

The absolute value is

$$
\begin{equation*}
\left|\frac{\left(V_{c d}\right)_{m}}{\left(E_{2}\right)_{m}}\right|=\frac{R_{s}}{\left(R_{0}+R_{s}\right) \omega R_{2}} \cdot \frac{G_{1} G_{2}}{m} \cdot \frac{1}{\sqrt{\left\{1-\left(T / T_{0}\right)^{2}\right\}^{2}+4 h^{2}\left(T / T_{0}\right)^{2}}} . \tag{27.1}
\end{equation*}
$$

The phase lag is

$$
\begin{equation*}
\operatorname{Arg}\left(\frac{\left(V_{c d}\right)_{m}}{x_{m}}\right)=-\frac{\pi}{2}-\tan ^{-1}\left[-2 h\left(T / T_{0}\right) / 1-\left(T / T_{0}\right)^{2}\right]+2 N \pi \tag{27.2}
\end{equation*}
$$

Define two functions $F_{3}(T)$ and $P_{3}(T)$ as follows.

$$
\begin{equation*}
F_{3}(T)=\frac{\left(R_{0}+R_{s}\right) \omega R_{2}}{R_{s}} \cdot\left|\frac{\left(V_{c d}\right)_{m}}{\left(E_{2}\right)_{m}}\right|=\frac{G_{1} G_{2}}{m} \cdot \frac{1}{\sqrt{\left\{1-\left(T / T_{0}\right)^{2}\right\}^{2}+4 h^{2}\left(T / T_{0}\right)^{2}}} \tag{28.1}
\end{equation*}
$$

where $\omega=2 \pi / T$. The high frequency asymptote of $F_{3}(T)$ gives the value of $G_{1} G_{2} / m$.

$$
\begin{equation*}
P_{3}(T)=\operatorname{Arg}\left(\frac{\left(V_{c d}\right)_{m}}{x_{m}}\right)+\frac{\pi}{2}=-\tan ^{-1}\left[-2 h\left(T / T_{0}\right) / 1-\left(T / T_{0}\right)^{2}\right]+2 N \pi \tag{28.2}
\end{equation*}
$$

$P_{3}(T)$ is equal to $\pi / 2$ at $T=T_{0}$
Note that $F_{3}(T)$ is almost same as $F_{1}(T)$ except of the value for high frequency asymptote, and $P_{3}(T)$ is the same function as $P_{1}(T)$. Then, the shape of $F_{3}(T)$ and $P_{3}(T)$ also can be seen in Fig.3.

Sensitivity of seismometer is defined as follows

$$
\begin{equation*}
G=l B \tag{29}
\end{equation*}
$$

where $l$ represents the length of coil wire within the magnetic field, $B$ denotes the flux density of magnetic field. Then, the ratio of the electro-dynamical constants of two coils can be given by the ratio of their resistances.

$$
G_{1} / G_{2}=l_{1} B / l_{2} B=l_{1} / l_{2}=R_{0} / R_{2}
$$

If we have the value of $m$, we can estimate the value of $G_{1}$ and $G_{2}$ independently.

$$
\begin{equation*}
G_{1}=\sqrt{\left(R_{0} / R_{2}\right) \cdot m F_{3}(0)} \cdot G_{2}=\sqrt{\left(R_{2} / R_{0}\right) \cdot m F_{3}(0)} . \tag{30}
\end{equation*}
$$

### 2.3.2 Step Current Input

Assume that the input voltage imbalance is a step function $E_{2}(t)=E_{0} \cdot H(t)$, where $H(t)$ denotes the Heviside function. The time domain solution of the equation of motion

$$
\ddot{x}(t)+2 \omega_{0} h \dot{x}(t)+\omega_{0}{ }^{2} x(t)=\frac{-G_{2} E_{0} H(t)}{m R_{2}},
$$

is given as

$$
\dot{x}(t)=\left\{\begin{array}{lc}
-G_{2} E_{0} / m R_{2} \cdot e^{-\omega_{0} h t} / \omega_{0} \sqrt{1-h^{2}} \cdot \sin \left(\omega_{0} \sqrt{1-h^{2}} \cdot t\right), h<1.0,  \tag{31}\\
-G_{2} E_{0} / m R_{2} \cdot t \cdot e^{-\omega_{0} t}, & h=1.0, \\
-G_{2} E_{0} / m R_{2} \cdot \sinh \left(\omega_{0} \sqrt{h^{2}-1} \cdot t\right), & h>1.0
\end{array}\right.
$$

The natural period and the damping constant can be estimated in the way described above Consideration on the circuit which consists of the seismometer and the shunt resistance $R_{s}$ gives

$$
\begin{equation*}
V_{c d}(t)=-\frac{G_{1}}{1+R_{0} / R_{S}} \cdot \dot{x}(t), \tag{32}
\end{equation*}
$$

Then,

$$
V_{a b}(t)=\left\{\begin{array}{lr}
\left(G_{1} G_{2} E_{0}\right) /\left[m R_{2}\left(1+R_{0} / R_{s}\right)\right] \cdot e^{-\omega_{0} h t} / \omega_{0} \sqrt{1-h^{2}} \cdot \sin \left(\omega_{0} \sqrt{1-h^{2}} \cdot t\right), & h<1.0, \\
\left(G_{1} G_{2} E_{0}\right) /\left[m R_{2}\left(1+R_{0} / R_{s}\right)\right] \cdot t \cdot e^{-\omega_{0} t}, & h=1.0, \\
\left(G_{1} G_{2} E_{0}\right) /\left[m R_{2}\left(1+R_{0} / R_{s}\right)\right] \cdot \sinh \left(\omega_{0} \sqrt{h^{2}-1} \cdot t\right), & h>1.0 .
\end{array}\right.
$$

The amplitude of the first peak is given

$$
\begin{equation*}
\left(V_{c d}\right)_{f . p .}=\frac{G_{1} G_{2} E_{0}}{m R_{2}\left(1+R_{0} / R_{s}\right) \omega_{0} P(h)}, \tag{33}
\end{equation*}
$$

where the suffix $f$.p. denotes the amplitude of the first peak, $p(h)$ is given by Eq.(20).
The value of $G_{l} G_{2} / m$ is given in the following,

$$
\begin{equation*}
\frac{G_{1} G_{2}}{m}=\frac{\left|\left(V_{c d}\right)_{f \cdot p \cdot \mid}\right| R_{2}\left(1+R_{0} / R_{s}\right) \omega_{0} P(h)}{E_{0}} . \tag{34}
\end{equation*}
$$

## 3. Examples of Application

### 3.1 Experiments on Shaking Table

The vertical component of the seismometer L-22D manufactured by Mark Product Ltd. is selected for the object of experiment. This is so old that its dynamic characteristics have been changed already. Moreover, the housing can be opened easily. These are enough for us to have doubt that the values of parameter in the specification sheet are unreliable.

Electro - Dynamic Test System in the Structural Testing Laboratory of the BRI consists of Function generator, Controller, Band Pass Filter, Shaker which gives the input motion to the tested seismometer, two Shaking tables, Vibrometer and Hydraulic unit. Shaker and Shaking table are located on the platform isolated from the rest of the building (Fig.7, Akashi, 1974). One of the shaking tables corresponds to vertical motion and another horizontal one. Each shaking table has a reference seismometer which is the differential transformer. The objective seismometer is fixed on the shaking table and its output is conserved on the recording paper simultaneously with the output of the reference seismometer.


Fig. 7 Electro-Dynamic Shaking Test System

At first, it is evaluated that the sensitivity of the reference seismometer is approximately 200 Volt per meter. The calibration process is the followings. Sinusoidal motion is provided to the shaking table by Controller and the amplitude of table's displacement relative to the flame is measured by one's eye with a scale. This amplitude is compared with the amplitude of the output voltage change of the reference seismometer. Thus, high accuracy can not be expected. At least an error of $5 \%$ must be taken into account for the sensitivity of the reference seismometer.

Two sinusoidal curves are drawn on the recording paper which is the output voltage of the tested seismometer and that of the reference one. Besides, the timing signal is drawn by the inner clock of the recorder. Then, we can determine the period or the frequency of the signals, the amplitudes of two seismograms and the delay of the peak between two sinusoidal records (Fig.8).

The period $T$ in second was calculated by the following.

$$
\begin{equation*}
T=D_{s} / D_{c}, \tag{42}
\end{equation*}
$$

where $D_{s,}$ is the length of one cycle of the signal on the recording paper measured in meter, $D_{c}$ the length measured in meter on the recording paper which corresponds to one second.

A similar procedure is applied to find the amplitude voltage values of the tested and the reference seismometers. The displacement amplitude $y_{m}(f)$ measured in meter of the shaking table is given by the following.

$$
\begin{equation*}
y_{m}(f)=\left\{A_{r e f}(f) \square G_{r e c}\right\} / G_{r e f} \tag{43}
\end{equation*}
$$

where $A_{\text {ref }}(f)$ measured in meter is the amplitude of the output the reference seismometer on the recording paper, $G_{\text {rec }}$, the gain of the recorder measured Volt per meter, $G_{\text {ref }}$ the sensitivity of the reference seismometer which is 200 Volt per meter in this case. The amplitude of the output voltage of the tested seismometer $e_{m}(f)$ measured in Volt e given as,

$$
\begin{equation*}
e_{m}(f)=A(f) \cdot G_{\text {rec }}{ }^{v}, \tag{44}
\end{equation*}
$$

where $A(f)$ measured in meter is the amplitude of the output of the tested seismometer on the recording paper, the gain of the recorder measured in Volt per meter.

The function $P_{l}(f)$ measured in degree is given by the following.
$P_{1}(f)=360-\left\{L(f) / D_{s} \cdot 360\right\}+90$.
where $L(f)$ is the delay measured in meter on the recording paper which is the length from the reference seismogram's peak to tested one's peak.

The measurements are conducted for five values of shunt resistance. Calculated $F_{l}(T)$ and $P_{l}(T)$ are plotted in Fig. 9 in comparison with several theoretical curves. The value of $G_{l}(70 \mathrm{v} / \mathrm{m} / \mathrm{sec})$ is estimated from the high frequency asymptote of $F_{l}(T)$ and the natural period $T_{0}(0.42 \mathrm{sec})$ from the intercept of $P_{l}(T)$ with 90 degree. Values of the damping constant are evaluated by comparison with theoretical curves. The correspondence of the damping constant to the shunt resistance is 0.38 -infinite, $0.57-10 \mathrm{Kolm}, 0.92-2 \mathrm{Kohm}$, 1.15-1 Kohm and 1.45-200 Ohm.


Fig. 8 Measurement on recording paper for sinusoidal input


Fig. 9 Calculated $F_{l}(T)$ and $P_{l}(T)$ plotted in comparison with several theoretical curves.

The damping constant $h$ is given in the form

$$
\begin{equation*}
h=h_{m}+\frac{G^{2}}{2 m \omega_{0}\left(R_{0}+R_{s}\right)}, \tag{46}
\end{equation*}
$$

where $m$ is the mass of the pendulum. Fig. 10 shows the plot of $h$ against $l /\left(R_{0}+R_{s}\right)$ for the data obtained above, where $R_{0}$ is equal to 2.25 Kohm for this experiment. The line in Fig. 10 shows a good fitting to the plotted data, of which formula is

$$
\begin{equation*}
h=0.38+2.4 \times 10^{3} /\left(R_{0}+R_{s}\right) . \tag{47}
\end{equation*}
$$

where two resistance are measured in ohm.
The slope of the line gives

$$
\begin{equation*}
G^{2} / m=2.4 \times 10^{3} \cdot 4 \pi / T_{0}=7.18 \times 10^{4}\left(\mathrm{v}^{2} / \mathrm{Kg} \mathrm{~m}^{2} / \mathrm{sec}^{2}\right) . \tag{48}
\end{equation*}
$$

This and the value of $G$ evaluated above give that of pendulum mass $m$

$$
m=0.068(\mathrm{~kg})=68(\mathrm{~g})
$$

The value of the shunt resistance which corresponds to $h=0.64$ can be evaluated by Fig.10. $R_{s}=7.45 \mathrm{Kohm}$ may give an appropriate value of the shunt resistance.

The estimated value of mass $m=0.068(\mathrm{Kg})$ slightly differs from that given in the specification sheet $m=0.0728(\mathrm{Kg})$. In this experiment, the sensitivity of the reference seismometer is assumed to be $200(\mathrm{v} / \mathrm{m})$, and can have an error of, e.g., about $5 \%$. If the sensitivity of the reference seismometer is assumed to be $193.6(\mathrm{v} / \mathrm{m})$, the $G$ obtained from the high frequency asymptote of $F_{1}(f)$ will be $72.3(\mathrm{v} / \mathrm{m} / \mathrm{sec})$ and the mass $m$ is $0.0728(\mathrm{Kg})$ for the same measured data. Thus, this discrepancy of the value for $m$ may be caused by the rough estimation of the sensitivity of the reference seismometer.


Fig. $10 h$ against $l /\left(R_{0}+R_{s}\right)$ obtained in the calibration on Shaking Table

### 3.2. Experiment using Willmore Impedance Bridge

The same L-22D seismometer used above is the object of the experiment. Similar process of measurement on the recording paper as described above are conducted (See Fig.8) with impedance bridge ( $R=100 \mathrm{Kohm}$ ).

### 3.2.1. Sinusoidal Current Input

Then, $F_{2}(T)$ and $P_{2}(T)$ are obtained as shown in Fig.11. Fig. 12 shows $h$ against $l /\left(R_{0}+R_{s}\right)$ for the data obtained above, where $R_{0}$ is equal to 2.25 Kohm for this experiment. The values of $h$ and $R_{s}$ are ( 0.38 -infinitive, $0.57-10$ Kohm, $0.71-5.1$ Kohm, $0.92-2$ Kohm, $1.15-1 \mathrm{Kohm}, 1.45-200 \mathrm{Ohm}$ ). The line in Fig. 12 is

$$
h=0.38+2.45 \times 10^{3} /\left(R_{0}+R_{s}\right) .
$$

This is consistent for the values $\left.G^{2} / m=7.0 \times 10^{4}\left(\mathrm{v}^{2} / \mathrm{Kg} \mathrm{m}^{2} / \mathrm{sec}^{2}\right)\right), T_{0}=0.42(\mathrm{sec})$ obtained from Fig. 11 .
The value of the shunt resistance which corresponds to $h=0.64$ can be evaluated by Fig. 12 That is $R_{s}=7.18$ Kohm which is the appropriate value of the shunt resistance.


Fig. $11 F_{2}(T)$ and $P_{2}(T)$ in comparison with several theoretical curves.


Fig. $12 h$ against $l /\left(R_{0}+R_{s}\right)$ for the result shown in Fig. 11.


Fig. 13 Measurement on recording paper for step current input

### 3.2.2. Step Current Input

Two wave trains are drawn on the recording paper which are $V_{a b}$ and $V_{A B}=E_{0}$. Beside of them, the timing signal is drawn by the inner clock of the recorder.

The damping ratio $v$ and damping factor $h$ were calculated by the following formula.

$$
v=a_{1} / a_{2}=a_{2} / a_{3}=a_{3} / a_{4}, h=0.733 \log _{10} v
$$

where $a_{1}$ is amplitude of the lust peak of the output voltage of the tested seismometer, $a_{2}$ is that of the second peak, etc.. For the overdamped case and the case $h$ near to 1.0, the estimation used $T_{w}$ is conducted. The time in that $E(t)$ takes the amplitude which is 0.2 times of the first peak $T_{w}$ is also measured on the recording paper (Fig.13).

The apparent period $T_{d}(\mathrm{sec})$ is calculated by the following formula with the measured values of $D_{s}$ and $D_{c}$.

$$
T_{d}=D_{s} / D_{c}
$$

where $D_{s}$ is the length in meter of one cycle of the signal on the recording paper, $D_{c}$ the length in meter on the recording paper which corresponds to one second. The natural period $T_{0}$ in second and the natural angular frequency $\omega_{0}$ is calculated with the estimated values of $h$ and $T_{d}\left(T_{0}=\sqrt{1-h^{2}} \cdot T_{d,}, \omega_{0}=2 \pi / T_{0}\right)$.

The amplitude of the input voltage of the tested seismometer $E_{0}(f)$ is given ,

$$
\begin{equation*}
E_{0}=A_{0} \cdot G_{r e c}, \tag{49}
\end{equation*}
$$

where $A_{0}$ is the amplitude of the step of input function measured in meter on the recording paper, $G_{\text {rec }}$ the gain of the recorder in volt per meter. The amplitude of the output voltage of the tested seismometer $\left(V_{a b}\right)_{f p}$. is given,

$$
\begin{equation*}
\left(V_{a b}\right)_{f p}=A \cdot G_{\text {rec }}{ }^{v}, \tag{50}
\end{equation*}
$$

where $A$ is the first peak's amplitude of the output of the tested seismometer on the recording paper, $G_{\text {rec }}^{v}$ the gain of the recorder.

The obtained values of $h$ and corresponding $R_{s}$ are ( 0.38 -infinite, $0.57-10 \mathrm{Kohm}, 0.71-5.1 \mathrm{Kohm}$, 0.92-2 Kohm, $1.15-1 \mathrm{Kohm}, 1.45-200 \mathrm{Ohm})$. The calculated value of the natural period is $T_{0}=0.42$ second and the value of $G^{2} / m$ is equal to $7.2 \times 10^{4}\left(\mathrm{v}^{2} / \mathrm{Kg} \mathrm{m}{ }^{2} / \mathrm{sec}^{2}\right)$.
$h$ against $l /\left(R_{0}+R_{s}\right)$ for the obtained data is shown in Fig.14, where $R_{0}$ is equal to 2.25 Kohm for this experiment. The line in Fig. 14 is

$$
h=0.38+2.45 \times 10^{3} /\left(R_{0}+R_{s}\right) .
$$

This is consistent for the $G^{2} / m$ obtained above from peak amplitude. The value of the shunt resistance which corresponds to $h=0.64$ can be evaluated by Fig.14. An appropriate value of the shunt resistance is given by $R_{s}=7.18$ Kohm.

The sensitivity calculated by $G^{2} / m=7.2 \times 10^{4}\left(\mathrm{v}^{2} / \mathrm{Kg}\right.$ $\left.\mathrm{m}^{2} / \mathrm{sec}^{2}\right)$ and $m=0.0728(\mathrm{Kg})$ is $G=72.4(\mathrm{~V} / \mathrm{m} / \mathrm{sec})$. This is acceptable
 compared with the value in the specification sheet $\left(G=75.7 \times 10^{4}\right)$.

### 3.3. Experiment using Calibration Coil

The vertical component of the seismometer L-4C-3D manufactured by Mark Product Ltd. is selected for the subject of practice. The weight of the pendulum mass is $0.96(\mathrm{Kg})$ according to the specification sheet. The transducer coil resistance is 5.5 Kohm , where as the calibration coil resistance is 19 ohm .

### 3.3.1. Sinusoidal Current Input

Similar measurement on recording paper described in Fig. 8 is conducted at various frequencies between 0.5 and 30.0 Hz with infinite value of shunt resistance. The results are plotted in Fig. 15 with the theoretical curve of which parameters are $G_{l} G_{2} / m=124.16\left(\mathrm{v}^{2} / \mathrm{Kg} \mathrm{m} / \mathrm{sec}^{2}\right), m=0.96 \mathrm{Kg}$., $T_{0}=0.935$ second and $h=0.23$.

These give
$G_{I}=185.7(\mathrm{v} / \mathrm{m} / \mathrm{sec}), G_{2}=0.642(\mathrm{v} / \mathrm{m} / \mathrm{sec})$.


Fig. $15 F_{3}(T)$ and $P_{3}(T)$ in comparison with theoretical curve

### 3.3.2. Step Current Input

Similar measurement as described in Fig. 13 is conducted for nine values of shunt resistance. The obtained values of $h$ and corresponding $R_{s}$ are ( 0.23 -infinitive, 0.29-100 Kohm, 0.37-51 Kohm, 0.52-20 Kohm 0.63-10 Kohm, $0.76-5.1 \mathrm{Kohm}, 1.02-2 \mathrm{Kohm}, 1.1-1 \mathrm{Kolm}, 1.24-510 \mathrm{Ohm}$ ). The natural period is $T_{0}=0.935$ second and the value of $G_{l} G_{2} / m$ is equal to $124.14\left(\mathrm{v}^{2} / \mathrm{Kg} \mathrm{m} / \mathrm{Kec}^{2}\right) . h$ against $l /\left(R_{0}+R_{s}\right)$ for the data obtained above is plotted in Fig.16, where $R_{0}$ is equal to 5.5 Kohm for this experiment. The curve in Fig. 16 is

$$
h=0.23+5.95 \mathrm{X10}^{3} /\left(R_{0}+R_{s}\right)
$$

This is consistent for the values $G_{l} G_{2} / m$ evaluated above. By using the value of mass in the specification sheet and those of coil resistances, the following two parameters are separated.

$$
G_{I}=185.74(\mathrm{v} / \mathrm{m} / \mathrm{sec}), G_{2}=0.642(\mathrm{v} / \mathrm{m} / \mathrm{sec})
$$



Fig16 $h$ against $l /\left(R_{0}+R_{s}\right)$

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